# Attractors in affine differential systems with impulsive control \*

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**Abstract.** In this paper we prove that an asymptotic equilibrium of an affine system of differential equations can become a strange attractor under affine impulsive control. The linear oscillator is studied as example.

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## 1 Introduction

We are concerned with the system of affine differential equations

$$\dot{x} = Ax + b \quad (x \in \mathbb{R}^m),\tag{1}$$

where A is a nonsingular matrix.

Suppose that at the moments  $t = n \in \mathbb{N}^*$  instantaneous control actions occur, which change the state of system as follows:

$$\Delta x \big|_{t=n} := x(n+0) - x(n-0) = C_{i_n} x(n-0) + d_{i_n} \quad (n \in \mathbb{N}^*),$$
(2)

where the matrices  $C_{i_n}$  and vectors  $d_{i_n}$  belong to given sets (finite or infinite).

Between any two consecutive kicks the motion of the system obeys (1). At the moment t = n the elements  $C_{i_n}$  and  $d_{i_n}$ , which determine the jump by (2), are chosen, say randomly. For convenience, we will consider that all solutions of system (1)-(2) are right continuous at the moments  $t \in \mathbb{N}$ .

Let  $\mathcal{F}$  be the set of all affine maps  $\{F_{i_n} : x \mapsto C_{i_n}x + d_{i_n}\}_{n \in \mathbb{N}^*}$ . For simplicity, we assume that the set  $\mathcal{F}$  is finite and contains only r distinct elements, say  $\mathcal{F} = \{F_1, F_2, \ldots, F_r\}$ . Denote  $\tilde{F}_n = E + F_n$   $(1 \le n \le r)$ , where E is the identity operator.

Assume that the spectra of the operators  $(E + C_n)e^A$   $(1 \le n \le r)$  are located strictly inside of the unit circle. Sometimes, if necessary, it is required that all operators  $E + C_n$   $(1 \le n \le r)$  are invertible.

It is known (see, e.g., [1]) that if the sequence  $\{F_{i_n}\}_{n\in\mathbb{N}^*}$  is periodic, then the behavior of the system is quite simple: there exists a globally attracting periodic cycle, corresponding to a periodic motion. This situation may, however, be changed essentially in the general case.

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In what follows we show that there exists a global attractor in the extended phase space and for "typical" sequence  $\{F_{i_n}\}_{n \in \mathbb{N}^*}$  this invariant set has a non-integer Hausdorff dimension (see, e.g., [2]) and represents a fractal. Moreover, the motions on the attractor are chaotic by Li-Yorke (in the meaning that every trajectory on the attractor is dense) and, as a consequence, every solution of the impulsive system (1)-(2) tends to a chaotic one.

This paper represents an application of general results from [3, 4].

## 2 Invariant sets

If one denotes by x(n) the state of the system immediately after the *n*-th kick, then by a straightforward calculation we end up with a sequence of affine maps  $\Phi_{i_n}: x(n) \mapsto x(n+1), n \in \mathbb{N}^*$ . More precisely,

$$\Phi_{i_n} : x \mapsto (E + C_{i_{n+1}})e^A x + (E + C_{i_{n+1}})(e^A - E)A^{-1}b + d_{i_{n+1}}$$

We call the sequence  $\{\Phi_{i_n}\}_{n\in\mathbb{N}}$  the *Poincaré system* associated with the impulsive system (1)–(2). Under the above assumptions, this Poincaré system is generated by r affine maps  $\{\Phi_1, \Phi_2, \ldots, \Phi_r\}$ .

Let  $\mathcal{P}_{cp}(\mathbb{R}^m)$  be the set of all nonempty compact subsets of  $\mathbb{R}^m$ , endowed with the Hausdorff-Pompeiu metrics (see, e.g.,[2]). Denote by  $\Phi$  the *Nadler-Hutchinson* operator [6] on the space  $\mathcal{P}_{cp}(\mathbb{R}^m)$ , defined by  $\Phi(M) := \bigcup_{n=1}^r \Phi_n[M], M \in \mathcal{P}_{cp}(\mathbb{R}^m)$ .

Let  $\psi(\cdot, \tau, x_0)$  stand for the solution of the system (1)–(2) with the initial condition  $x(\tau) = x_0$ . Since the system (1)–(2) is affine, there is a unique such solution, defined on  $\mathbb{R}$  (see, e.g., [1]).

**Lemma 1.** For every  $x \in \mathbb{R}^m$  and  $t_1, t_2, t_3 \in \mathbb{R}$  one has

$$\psi(t_1, t_2, \psi(t_2, t_3, x)) = \psi(t_1, t_3, x).$$
(3)

**Proof.** This follows immediately from the uniqueness of the solution of respective Cauchy problem.  $\hfill \Box$ 

**Lemma 2.** Every solution of the system (1)-(2) can be written as follows:

$$\begin{split} \psi(t,\tau,x) &= e^{(t-\tau)A}x, \quad if \ [t] = [\tau], \quad or \\ \psi(t,\tau,x) &= e^{(t-[t])A}\tilde{F}_{i_{[t]}}e^{([t]-\tau)A}x, \quad if \ [t] = [\tau]+1, \quad or \\ \psi(t,\tau,x) &= e^{(t-[t])A}\left(\prod_{j=[\tau]+2}^{[t]} \Phi_{i_j}\right)\tilde{F}_{i_{[\tau]+1}}e^{([\tau]+1-\tau)A}x, \quad if \ [t] > [\tau]+1, \quad or \\ \psi(t,\tau,x) &= e^{(t-[\tau])A}\tilde{F}_{i_{[\tau]}}^{-1}e^{([\tau]-\tau)A}x, \quad if \ [t] = [\tau]-1, \quad or \\ \psi(t,\tau,x) &= e^{(t-[t]-1)A}\tilde{F}_{i_{[\tau]}+1}^{-1}\left(\prod_{j=[\tau]}^{[t]+2} \Phi_{i_j}^{-1}\right)e^{([\tau]-\tau)A}x, \quad if \ [t] < [\tau]-1, \end{split}$$

where  $[\cdot]$  denotes the integral part.

**Proof.** The proof is straightforward.

**Remark 1.** In fact, the impulsive system (1)–(2) is nonautonomous and we can consider only its integral curves. Even in the case of periodic sequence  $\{F_{i_n}\}_{n \in \mathbb{N}^*}$  we cannot factorize the system to obtain a (autonomous) system on the direct product  $S^1 \times \mathbb{R}^m$ . However, since the impulse actions occur at the moments  $t \in \mathbb{N}$ , we can obtain a foliation on the cylinder  $S^1 \times \mathbb{R}^m$  by factorization on time. This foliation consists of pieces of integral curves of the system (1)–(2).

Project the system (1)-(2) to the cylinder  $S^1 \times \mathbb{R}^m$ , using the projection  $\pi$ :  $(t, y) \mapsto (t \pmod{1}, y).$ 

We will say that a set  $V \subset S^1 \times \mathbb{R}^m$  is *positive invariant* (*invariant*) with respect to the system (1)-(2), if for every point  $(\tau, x) \in V$  and every natural k one has

$$(t \pmod{1}, \psi(t, \tau + k, x)) \in V \text{ for } t \ge \tau + k \pmod{t \in \mathbb{R}}.$$
(4)

In other words, V consists of pieces of integral curves.

By definition, such a set  $V \subset S^1 \times \mathbb{R}^m$  covers the whole base  $S^1$  by projection. Denote by  $(t, V_t) := \{(t, x) \in S^1 \times \mathbb{R}^m | (t, x) \in V\}$  the *fiber* over the point  $t \in [0, 1)$ . For convenience, we will identify this fiber with  $V_t$ . Moreover, in the sequel the notation  $V_t$  for  $t \in \mathbb{R}$  will mean  $V_t \pmod{1}$ .

**Theorem 3.** The set  $V \subset S^1 \times \mathbb{R}^m$  is positive invariant (invariant) with respect to the system (1)-(2) if and only if it satisfies the following conditions:

1. 
$$e^{(t-\tau)A}V_{\tau} \subset V_t$$
 for  $0 \le \tau \le t < 1$   $(e^{(t-\tau)A}V_{\tau} = V_t$  for  $\tau, t \in [0,1)$ );  
2.  $\bigcup_{n=1}^r \tilde{F}_n e^{(1-\tau)A}V_{\tau} \subset V_0$   $(\bigcup_{n=1}^r \tilde{F}_n e^{(1-\tau)A}V_{\tau} = V_0)$  for  $0 \le \tau < 1$ .

**Proof.** Necessity. Assume that the set  $V \subset S^1 \times \mathbb{R}^m$  is positive invariant (invariant). If  $x \in V_{\tau}$ , then  $(\tau, x) \in V$ , and by Lemma 2 one has for  $0 \leq \tau \leq t < 1$ 

$$e^{(t-\tau)A}x = \psi(t,\tau,x) \in V_t.$$
(5)

At the same time, for every natural n there is a natural  $k_n$  such that

$$\tilde{F}_n e^{(1-\tau)A} x = \psi([\tau] + k_n + 1, \tau + k_n, x) \in V_{[t]+k_n+1} = V_0.$$

In the case of invariance the relation (5) holds for  $\tau, t \in [0, 1)$ . Moreover, for every  $z \in V_t$  there is  $y = \psi(\tau, t, z) \in V_\tau$  which verifies  $e^{(t-\tau)A}y = z$ .

Analogously, for every  $0 \leq \tau < 1$ ,  $z \in V_0$  and  $1 \leq n \leq r$  there is a natural  $k_n$  such that  $\tilde{F}_{k_n} = \tilde{F}_n$  and there is  $y = \psi(\tau + k_n - 1, k_n, z) \in V_{\tau+k_n-1} = V_{\tau}$ , verifying  $\tilde{F}_n e^{(1-\tau)A} y = \psi(k_n, \tau + k_n - 1, y) = z \in V_0$ .

Sufficiency. Assume that conditions 1)–2) hold. Let  $(\tau, x) \in V$ . To proof (5) take firstly  $t \ge \tau + k, k \in \mathbb{N}^*$ .

If  $[t] = [\tau] + k$  or  $[t] = [\tau] + k + 1$ , then (4) is a consequence of the conditions 1)–2) and Lemma 2.

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If  $[t] = q, [\tau] + k = p, q > p + 1$ , then by Lemmas 1 and 2 the conditions 1)–2) imply for every  $x \in V_{\tau}$ :

$$\psi(t,\tau+k,x) = e^{(t-q)A} \prod_{j=p+2}^{q} \Phi_{i_j} \tilde{F}_{i_{p+1}} e^{(p+1-\tau-k)A} x \in V_t,$$
  
$$\psi(t,\tau+k,x) = \psi(t,q,\psi(q,p+1,\psi(p+1,\tau+k,x))) \in V_t.$$

In the case of invariance we consider in addition  $t < [\tau] + k$ . In this case there is  $y \in V_0$  such that  $x \in V_{\tau} = V_{\tau+k}$  may be represented as  $x = \psi(\tau + k, [\tau] + k, y)$ . In turn, there is  $z \in V_{[t]} = V_0$  such that  $y = \psi([\tau] + k, [t], z)$ . By Lemma 1,  $\psi(t, \tau + k, x) = \psi(t, [\tau] + k, y) = \psi(t, [t], z) \in V_t$ . This completes the proof.  $\Box$ 

**Corollary 4.** If  $V \subset S^1 \times \mathbb{R}^m$  is positive invariant (invariant), then  $V_0$  is positive invariant (invariant) with respect to the Nadler-Hutchinson operator  $\Phi$ , i.e.  $\Phi(V_0) = \bigcup_{n=1}^r \Phi_n[V_0] \subset V_0$  ( $\Phi(V_0) = V_0$ ).

## **3** IFS and attractors

Since the eigenvalues of all matrices  $(E + C_n)e^A$   $(1 \le n \le r)$  are located strictly inside the unit circle, all operators  $\Phi_n$   $(1 \le n \le r)$  are contracting.

We associate to the system (1)-(2) a hyperbolic Iterated Function System (IFS)  $\{\mathbb{R}^m; \Phi_1, \Phi_2, \ldots, \Phi_r\}$  (see, e.g., [2]), consisting of affine contractions. This IFS determines in  $\mathbb{R}^m$  a global compact attractor K, which is the unique fixed point of the corresponding contractive Nadler-Hutchinson operator  $\Phi$ .

Given the natural k we say that the sequence  $\{F_{i_n}\}_{n\in\mathbb{N}^*}$  is k-universal if it contains every word of the length k from the alphabet of  $\mathcal{F} = \{F_1, F_2, \ldots, F_r\}$ . We say that the sequence  $\{F_{i_n}\}_{n\in\mathbb{N}^*}$  is universal if it is k-universal for every natural k.

By Lemma Borel-Cantelli [5], if the sequence  $\{F_{i_n}\}_{n\in\mathbb{N}^*}$  is chosen randomly with a uniform distribution, then with probability 1 it is universal.

If the sequence  $\{F_{i_n}\}_{n \in \mathbb{N}}$  is universal, then the orbit of each point in K is dense on K (is chaotic by Li-Yorke).

Recall some notions (see, e.g., [2]). A set is called *totally disconnected* if for every its point the connected component, containing this point, is the point itself. A set S is called *perfect* if it is closed and every point  $p \in S$  is the limit of points  $q_n \in S \setminus \{p\}$ . A set is called a *Cantor set* if it is totally disconnected, perfect and compact.

**Theorem 5.** If the spectra of the operators  $(E + C_n)e^A$   $(1 \le n \le r)$  are located strictly inside the disk of radius  $\frac{1}{r}$ , then the attractor K is totally disconnected.

**Proof.** It is known [7] that if an IFS consists of r contractions, each of them with the contraction coefficient s, and rs < 1, then the attractor K of this IFS is totally disconnected. It is sufficient to say that in our case  $s = \max_{1 \le n \le r} ||(E + C_n)e^A|| < \frac{1}{r}$ .  $\Box$ 

**Remark 2.** The hypotheses of Theorem 5 are far from being also necessary conditions for the attractor K to be totally disconnected.

Denote the distance from the point  $x \in \mathbb{R}^m$  to the compact  $M \subset \mathbb{R}^m$  by  $\varrho(x, M) := \min\{d(x, y) \mid y \in M\}.$ 

A bounded subset  $V \subset S^1 \times \mathbb{R}^m$  is called an *attractor* of the system (1)-(2) if it is positive invariant and for every solution  $\psi(\cdot, \tau, x)$  one has  $\varrho(\psi(t, \tau, x), V_t) \to 0$  as  $t \to +\infty$ .

**Theorem 6.** There exists an attractor of the system (1)-(2)

$$K^* = \left\{ (t, e^{tA}x + (e^A - E)A^{-1}b) \, | \, t \in [0, 1), x \in K \right\} \subset S^1 \times \mathbb{R}^m$$

with the Hausdorff dimension  $DH(K^*)$ , verifying the inequalities:

$$1 < DH(K^*) \le 1 - \frac{\ln r}{\ln s},$$
 (6)

where s is the smallest radius of a disc centered at the origin of coordinates, which contains the spectra of the operators  $(E + C_n)e^A$   $(1 \le n \le r)$ .

**Proof.** By Theorem 3, the set  $K^*$  is positive invariant. Since K is compact, every fiber  $K_t = e^{tA}K + (e^A - E)A^{-1}b$   $(0 \le t < 1)$  is compact as well. The compact K attracts every compact M in the fiber  $0 \times \mathbb{R}^m$  under the actions of the Nadler-Hutchinson operator  $\Phi$ . As a result every fiber  $K_t = e^{\nu A}K + (e^A - E)A^{-1}b$ , where  $\nu = t \pmod{1}$ , attracts the image  $e^{(t-[t])A}\Phi^n e^{([t]+1-t)A}M$  as  $n \to +\infty$  uniformly on  $t \in \mathbb{R}$ .

It is known (see, e.g., [2]) that the Hausdorff dimension of K verifies the inequality  $DH(K) \leq -\frac{\ln r}{\ln s}$ . This implies that  $DH(K^*) = 1 + DH(K) \leq 1 - \frac{\ln r}{\ln s}$ .

**Theorem 7.** Let the spectra of the operators  $(E + C_n)e^A$   $(1 \le n \le r)$  be located strictly inside the disk of radius  $\frac{1}{r}$  and let the sequence  $\{F_{i_n}\}_{n\in\mathbb{N}}$  be universal. Then the attractor  $K^*$  is homeomorphic to the direct product of a half-open interval and a Cantor set, its Hausdorff dimension verifying the inequalities:  $1 < DH(K^*) < 2$ .

**Proof.** It is easily seen from Theorem 6 that the attractor  $K^*$  is homeomorphic to the direct product  $[0,1) \times K$ . Under the given hypothesis, the compact K is perfect and by Theorem 5 is totally disconnected, and, as a consequence, it is a Cantor set. In this case the inequalities (6) become:  $1 < DH(K^*) < 2$ .

**Lemma 8.** There exist  $L > 0, \gamma > 0$  such that for any  $y \in \mathbb{R}^m$  there exists  $x \in K$ , satisfying

$$\|\psi(t,0,y) - \psi(t,0,x)\| \le Le^{-\gamma t} \|y - x\| \quad (t \ge 0).$$

**Proof.** This follows immediately from Lemma 2 and Theorem 6.

**Theorem 9.** If the sequence  $\{F_{i_n}\}_{n \in \mathbb{N}}$  is universal, then every integral curve of the system (1)–(2), starting in  $K^*$ , is chaotic by Li-Yorke, i.e. is dense in  $K^*$ .

**Proof.** If the sequence  $\{F_{i_n}\}_{n\in\mathbb{N}}$  is universal, then for every solution  $\psi(\cdot, \tau, x)$  with  $(\tau, x) \in K^*$ , the sequence  $\{\psi(j, \tau, x)\}_{j\geq [\tau]+1}$  is the orbit of the point  $\psi([\tau]+1, \tau, x) \in K$  under the IFS  $\{\mathbb{R}^m; \Phi_1, \ldots, \Phi_r\}$ . Since this orbit is dense on K, the integral curve of the solution  $\psi(\cdot, \tau, x)$  is dense on  $K^*$ .  $\Box$ 

**Corollary 10.** If the sequence  $\{F_{i_n}\}_{n \in \mathbb{N}}$  is universal, then every integral curve of the system (1)-(2) is chaotic by Li-Yorke or tends to a chaotic one.

**Proof.** This follows from Lemma 8 and Theorem 9.

Remark 3. If the impulses occur only in some integer moments, i.e.

$$\Delta x \big|_{t=\tau_n} := x(\tau_n + 0) - x(\tau_n - 0) = C_{i_n} x(\tau_n - 0) + d_{i_n} \quad (\tau_n \in \mathbb{N}^*), \tag{7}$$

then the system (1),(7) may be considered as a particular case of the system (1)-(2) by supplementing the set  $\mathcal{F}$  with the null operator F = 0 for other integer moments.

**Remark 4.** Analogously, if the spectra of all operators  $(E + C_n)e^A$   $(1 \le n \le r)$  are located strictly outside the unit circle, we can say about the repeller of the system (1)-(2).

**Remark 5.** Many classical fractals may be represented as attractors of affine IFS on  $\mathbb{R}^2$ . Fig. 1 shows some of them as attractors K of impulsive differential equations on  $\mathbb{C}$ , for example:

- the Sierpinski triangle in  $\dot{z} = -z \cdot \ln 2, \Delta z|_{t=n} = i \exp \frac{2\pi n i}{3}$   $(n \in \mathbb{N});$
- the pentagasket in  $\dot{z} = -z \cdot \ln(\frac{3+\sqrt{5}}{2}), \Delta z|_{t=n} = i \exp \frac{2\pi ni}{5} \ (n \in \mathbb{N});$
- the hexagasket in  $\dot{z} = -z \cdot \ln 3$ ,  $\Delta z|_{t=n} = \exp \frac{\pi n i}{3}$   $(n \in \mathbb{N})$ .



Figure 1. Fractals: the Sierpinski triangle, pentagasket and hexagasket as attractors K of impulsive affine systems

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## 4 Linear oscillator

Let us consider, as an example, the linear oscillator with impulsive actions

$$\ddot{x} + c\dot{x} + kx = 0 \quad (c > 0, k > 0), \tag{8}$$

$$\Delta \dot{x}\big|_{t=n} := \dot{x}(n+0) - \dot{x}(n-0) = \xi_{i_n} \quad (n \ge 1).$$
(9)

Assume that the range of the sequence  $\{\xi_{i_n}\}_{n\geq 1}$  contains only r distinct elements.

We can reduce the equations (8)–(9) to an affine system of impulsive differential equations (1)–(2) in the phase space  $x_1 = x, x_2 = \dot{x}$ . In this case we obtain some analogues of the previous theorems.

**Theorem 11.** [8] There exists an attractor  $K^* \subset S^1 \times \mathbb{R}^2_{(x,\dot{x})}$  of the system (8)-(9) with the Hausdorff dimension

$$DH(K^*) = 1 + \frac{2\sqrt{2}\ln r}{\sqrt{2}c - \sqrt{c^2 - 4k + |c^2 - 4k|}}.$$

**Theorem 12.** [8] If

$$2\ln r < c < \frac{k}{\ln r} + \ln r$$

and the sequence  $\{\xi_{i_n}\}_{n\in\mathbb{N}}$  is universal, then the attractor  $K^*$  is homeomorphic to the direct product of a half-open interval and a Cantor set, its Hausdorff dimension verifying the inequalities:  $1 < DH(K^*) < 2$ .

Fig. 2 represents the respective attractors  $K \subset \mathbb{R}^2_{(x,\dot{x})}$  for distinct values of parameters for two impulsive differential equations (8)–(9): on the left for c = 5/2, k = 2 and r = 3 (K is totally disconnected), on the right for c = 1, k = 5/4 and r = 3 (K is connected).



Figure 2. Attractor K for: c = 5/2, k = 2, r = 3 (left) and c = 1, k = 5/4, r = 3 (right)

**Remark 6.** If the range of sequence  $\{\xi_{i_n}\}_{n\geq 1}$  is infinite but bounded, then the system (8)–(9) admits an attractor as well.

Fig. 3 represents the attractor K (of an infinite IFS) for an impulsive differential equation (8)–(9) with c = 2, k = 2, where the values  $\{\xi_{i_n}\}_{n \ge 1}$  are randomly chosen from [0, 1].

All calculations and graphic objects have been done using the Computer Algebra System *Mathematica*.



Figure 3. Attractor K for  $c = 2, k = 2, \xi \in [0, 1]$ 

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