Characteristic functions of Markovian random evolutions in \mathbb{R}^m

Alexander D. Kolesnik

Abstract. The recurrent and integral relations for characteristic functions of Markovian random evolution in \mathbb{R}^m , $m \geq 2$, are presented.

Mathematics subject classification: Primary 60K99. Secondary 62G30, 60K35, 60J60, 60H30.

Keywords and phrases: Random evolution, characteristic functions, convolutions, Volterra integral equation.

The particular models of random evolutions in various Euclidean spaces of lower dimensions were studied in [1-5]. In this note we announce the recent results on the characteristic functions for the most general m-dimensional random evolution.

The subject of our interests is the following stochastic motion. A particle starts its motion from the origin $x_1 = \cdots = x_m = 0$ of the space \mathbb{R}^m , $m \ge 2$ at time t = 0. The particle is endowed with constant, finite speed c. The initial direction is a random m-dimensional vector with uniform distribution on the unit m-sphere

$$S_1^m = \left\{ \mathbf{x} = (x_1, \dots, x_m) \in \mathbb{R}^m : x_1^2 + \dots + x_m^2 = 1 \right\}.$$

The particle changes direction at random instants which form a homogeneous Poisson process of rate $\lambda > 0$. At these moments it instantaneously takes on the new direction with uniform distribution on S_1^m , independently of its previous motion.

Let $\mathbf{X}(t) = (X_1(t), \dots, X_m(t))$ be the position of the particle at an arbitrary time t > 0. At first, we concentrate our attention on the conditional distributions

$$Pr\{\mathbf{X}(t) \in d\mathbf{x} \mid N(t) = n\} =$$

$$= Pr\{X_1(t) \in dx_1, \dots, X_m(t) \in dx_m \mid N(t) = n\}, \quad n \ge 1$$

where N(t) is the number of Poisson events that have occurred in the interval (0, t)and $d\mathbf{x} = dx_1 \dots dx_m$ is the infinitesimal volume in the space \mathbb{R}^m .

Consider the conditional characteristic functions:

$$H_n(t) = E\left\{e^{i(\boldsymbol{\alpha}, \mathbf{X}(t))} | N(t) = n\right\}, \qquad n \ge 1,$$
(1)

where $\boldsymbol{\alpha} = (\alpha_1, \ldots, \alpha_m) \in \mathbb{R}^m$ is the real *m*-dimensional vector of inversion parameters and $(\boldsymbol{\alpha}, \mathbf{X}(t))$ denotes the scalar (inner) product of the vectors $\boldsymbol{\alpha}$ and $\mathbf{X}(t)$.

[©] Alexander D. Kolesnik, 2006

Computing the expectation in (1) we obtain

$$H_{n}(t) = \frac{n!}{t^{n}} \int_{0}^{t} d\tau_{1} \int_{\tau_{1}}^{t} d\tau_{2} \dots \int_{\tau_{n-1}}^{t} d\tau_{n} \times \left\{ \prod_{j=1}^{n+1} \left[2^{(m-2)/2} \Gamma\left(\frac{m}{2}\right) \frac{J_{(m-2)/2}(c(\tau_{j}-\tau_{j-1})\|\boldsymbol{\alpha}\|)}{(c(\tau_{j}-\tau_{j-1})\|\boldsymbol{\alpha}\|)^{(m-2)/2}} \right] \right\}.$$
 (2)

For the particular cases m = 2 (planar motion) and m = 4 (four-dimensional motion) the conditional characteristic functions (2) were explicitly computed in [2] (see formula (18) therein), and in [1] (see formula (15) therein), respectively.

We introduce the function

$$\varphi(t) = 2^{(m-2)/2} \Gamma\left(\frac{m}{2}\right) \frac{J_{(m-2)/2}(ct \|\boldsymbol{\alpha}\|)}{(ct \|\boldsymbol{\alpha}\|)^{(m-2)/2}}, \qquad m \ge 2.$$
(3)

Then (2) can rewritten in the following form

$$H_n(t) = \frac{n!}{t^n} \int_0^t d\tau_1 \int_{\tau_1}^t d\tau_2 \dots \int_{\tau_{n-1}}^t d\tau_n \left\{ \prod_{j=1}^{n+1} \varphi(\tau_j - \tau_{j-1}) \right\}, \qquad n \ge 1.$$
(4)

Denote the integral factor in (4) as follows

$$\mathcal{I}_{n}(t) = \int_{0}^{t} d\tau_{1} \int_{\tau_{1}}^{t} d\tau_{2} \dots \int_{\tau_{n-1}}^{t} d\tau_{n} \left\{ \prod_{j=1}^{n+1} \varphi(\tau_{j} - \tau_{j-1}) \right\}, \qquad n \ge 1.$$
(5)

The following theorem states that, for different $n \ge 1$, the functions (5) are connected with each other by a convolution-type recurrent relation. **Theorem 1.** For any $n \ge 1$ the following recurrent relation holds

$$\mathcal{I}_n(t) = \int_0^t \varphi(t-\tau) \,\mathcal{I}_{n-1}(\tau) \,d\tau = \int_0^t \varphi(\tau) \,\mathcal{I}_{n-1}(t-\tau) \,d\tau, \qquad n \ge 1, \qquad (6)$$

where, by definition, $\mathcal{I}_0(x) = \varphi(x)$.

Note that formula (6) can be rewritten in the following convolution form

$$\mathcal{I}_n(t) = \varphi(t) * \mathcal{I}_{n-1}(t), \qquad n \ge 1.$$
(7)

Corollary 1.1. For any $n \ge 1$ the following relation holds

$$\mathcal{I}_n(t) = \left[\varphi(t)\right]^{*(n+1)}, \qquad n \ge 1, \tag{8}$$

where the symbol *(n + 1) means the (n + 1)-multiple convolution.

Application of the Laplace transform

$$\mathcal{L}[f(t)](s) = \int_0^\infty e^{-st} f(t) \, dt, \qquad \text{Re } s > 0,$$

to the equality (8) leads to the following important result. Corollary 1.2. For any $n \ge 1$ the Laplace transform of functions (5) has the form

$$\mathcal{L}\left[\mathcal{I}_{n}(t)\right](s) = \left(\mathcal{L}\left[\varphi(t)\right](s)\right)^{n+1}, \qquad n \ge 1.$$
(9)

These results show that the function $\varphi(t)$ given by (3) plays a key role in our analysis. The reason is that $\varphi(t)$ is exactly the characteristic function (Fourier transform) of the uniform distribution on the surface of the *m*-sphere S_{ct}^m of the radius *ct*.

From both the Theorem 1 and its corollaries we see that the conditional characteristic functions $H_n(t)$ and their Laplace transforms, in fact, are expressed in terms of function $\varphi(t)$. Formula (9) shows that the possibility of obtaining the explicit form of the conditional characteristic functions (4) entirely depends on whether the exact Laplace transform of the function $\varphi(t)$ and its inverse Laplace transform can be explicitly computed.

Our next result presents a general formula for the conditional characteristic functions $H_n(t)$ in terms of inverse Laplace transform.

Theorem 2. For any $n \ge 1$ and any t > 0 the conditional characteristic functions (4) are given by

$$H_n(t) = \frac{n!}{t^n} \mathcal{L}^{-1} \left[\left(\frac{1}{\sqrt{s^2 + (c \|\boldsymbol{\alpha}\|)^2}} F\left(\frac{1}{2}, \frac{m-2}{2}; \frac{m}{2}; \frac{(c \|\boldsymbol{\alpha}\|)^2}{s^2 + (c \|\boldsymbol{\alpha}\|)^2} \right) \right)^{n+1} \right] (t),$$
(10)

where \mathcal{L}^{-1} means the inverse Laplace transform and

$$F(\xi,\eta;\zeta;z) = {}_{2}F_{1}(\xi,\eta;\zeta;z) = \sum_{k=0}^{\infty} \frac{(\xi)_{k}(\eta)_{k}}{(\zeta)_{k}} \frac{z^{k}}{k!}$$

is the standard hypergeometric function.

In view of (4), the characteristic function of $\mathbf{X}(t)$, $t \ge 0$, is given by the uniformly converging series

$$H(t) = e^{-\lambda t} \sum_{n=0}^{\infty} \lambda^n \, \mathcal{I}_n(t).$$
(11)

The following theorem presents the integral equation for the function H(t). **Theorem 3.** The characteristic function H(t), $t \ge 0$, satisfies the following convolution-type Volterra integral equation of second kind with the kernel $e^{-\lambda t}\varphi(t)$:

$$H(t) = e^{-\lambda t}\varphi(t) + \lambda \int_0^t e^{-\lambda(t-\tau)}\varphi(t-\tau)H(\tau) \, d\tau, \qquad t \ge 0.$$
(12)

The integral equation (12) can be rewritten in the following convolution form

$$H(t) = e^{-\lambda t} \varphi(t) + \lambda \left[\left(e^{-\lambda t} \varphi(t) \right) * H(t) \right], \qquad t \ge 0.$$
(13)

From this we immediately obtain the general formula for the Laplace transform of the characteristic function H(t):

$$\mathcal{L}[H(t)](s) = \frac{\mathcal{L}[\varphi(t)](s+\lambda)}{1-\lambda \mathcal{L}[\varphi(t)](s+\lambda)}, \qquad \text{Re } s > 0.$$
(14)

The explicit form of (14) is

$$\mathcal{L}[H(t)](s) = \frac{F\left(\frac{1}{2}, \frac{m-2}{2}; \frac{m}{2}; \frac{(c\|\boldsymbol{\alpha}\|)^2}{(s+\lambda)^2 + (c\|\boldsymbol{\alpha}\|)^2}\right)}{\sqrt{(s+\lambda)^2 + (c\|\boldsymbol{\alpha}\|)^2} - \lambda F\left(\frac{1}{2}, \frac{m-2}{2}; \frac{m}{2}; \frac{(c\|\boldsymbol{\alpha}\|)^2}{(s+\lambda)^2 + (c\|\boldsymbol{\alpha}\|)^2}\right)}.$$
(15)

From (11) and (8) it follows that the solution of equation (13) has the form

$$H(t) = e^{-\lambda t} \sum_{n=0}^{\infty} \lambda^n \ [\varphi(t)]^{*(n+1)} \ .$$
(16)

One should emphasize that, although formula (16) gives a general form of the characteristic function H(t), the multiple convolutions of the function $\varphi(t)$ with itself can scarcely be explicitly evaluated for arbitrary dimension.

From (12) we can see that

$$H(t)|_{t=0} = 1, \qquad \left. \frac{\partial H(t)}{\partial t} \right|_{t=0} = 0,$$

and, therefore, the transition density $f(\mathbf{x}, t)$, $\mathbf{x} \in \mathbb{R}^m$, $t \ge 0$, of the process $\mathbf{X}(t)$ satisfies the following initial conditions

$$f(\mathbf{x},t)|_{t=0} = \delta(\mathbf{x}), \qquad \left. \frac{\partial f(\mathbf{x},t)}{\partial t} \right|_{t=0} = 0,$$

where $\delta(\mathbf{x})$ is the *m*-dimensional Dirac delta-function.

References

- [1] KOLESNIK A.D. A four-dimensional random motion at finite speed. J. Appl. Prob., 2006, 43.
- [2] KOLESNIK A.D., ORSINGHER E. A planar random motion with an infinite number of directions controlled by the damped wave equation. J. Appl. Prob., 2005, 42, p. 1168–1182.
- [3] MASOLIVER J., PORRÁ J.M., WEISS G.H. Some two and three-dimensional persistent random walks. Physica A, 1993, 193, p. 469–482.
- [4] STADJE W. The exact probability distribution of a two-dimensional random walk. J. Stat. Phys., 1987, 46, p. 207–216.
- [5] STADJE W. Exact probability distributions for non-correlated random walk models. J. Stat. Phys., 1989, 56, p. 415–435.

Institute of Mathematics and Computer Science Academy of Sciences of Moldova Academiei str. 5, MD-2028 Kishinev Moldova E-mail: kolesnik@math.md Received November 3, 2006