

On Riemann extension of the Schwarzschild metric

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Abstract. The properties of the Riemann extension of the Schwarzschild metric are studied.

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1 Introduction

The notice of the Riemann extension of nonriemannian spaces was first introduced in ([1]). Main idea of this theory is application of the methods of Riemann geometry for studying of the properties of nonriemannian spaces.

For example the system differential equations in form

$$\frac{d^2 x^k}{ds^2} + \Pi_{ij}^k \frac{dx^i}{ds} \frac{dx^j}{ds} = 0 \quad (1)$$

with arbitrary coefficients $\Pi_{ij}^k(x^l)$ can be considered as the system of geodesic equations of affinely connected space with local coordinates x^k .

For the n -dimensional Riemannian spaces with the metrics

$${}^n ds^2 = g_{ij} dx^i dx^j$$

the system of geodesic equations looks same but the coefficients $\Pi_{ij}^k(x^l)$ now have very special form and depends from the choice of the metric g_{ij} .

$$\Pi_{kl}^i = \Gamma_{kl}^i = \frac{1}{2} g^{im} (g_{mk,l} + g_{ml,k} - g_{kl,m})$$

In order that the methods of Riemann geometry can be applied for studying of the properties of the spaces with equations (1) the construction of $2n$ -dimensional extension of the space with local coordinates x^i was introduced .

The metric of extended space constructs with help of coefficients of equation (1) and looks as follows

$${}^{2n} ds^2 = -2\Pi_{ij}^k(x^l) \Psi_k dx^i dx^j + 2d\Psi_k dx^k \quad (2)$$

where Ψ_k are the coordinates of additional space.

The important property of such type metric is that the geodesic equations of metric (2) consist from the two parts

$$\ddot{x}^k + \Gamma_{ij}^k \dot{x}^i \dot{x}^j = 0, \quad (3)$$

and

$$\frac{\delta^2 \Psi_k}{ds^2} + R_{kji}^l \dot{x}^j \dot{x}^i \Psi_l = 0, \quad (4)$$

where

$$\frac{\delta \Psi_k}{ds} = \frac{d\Psi_k}{ds} - \Pi_{jk}^l \Psi_l \frac{dx^j}{ds}.$$

The first part (3) of complete system is the system of equations for geodesics of basic space with local coordinates x^i and they does not contains the coordinates Ψ_k .

The second part (4) of system of geodesic equations has the form of linear 4×4 matrix system of second order ODE's for coordinates Ψ_k

$$\frac{d^2 \vec{\Psi}}{ds^2} + A(s) \frac{d\vec{\Psi}}{ds} + B(s) \vec{\Psi} = 0. \quad (5)$$

From this point of view we get the case of geodesical extension of basic space in local coordinates (x^i) .

It is important to note that the geometry of extended space is connected with geometry of basic space. For example the property of this space to be Ricci-flat keeps also for the extended space.

This fact give us the possibility to use the linear system of equation (5) for studying of the properties of basic space.

In particular the invariants of the 4×4 matrix-function

$$E = B - \frac{1}{2} \frac{dA}{ds} - \frac{1}{4} A^2$$

under change of the coordinates Ψ_k can be used for that.

The first applications of the notice of extended spaces the studying of nonlinear second order differential equations connected with nonlinear dynamical systems was done in works of author ([2-4]).

Here we consider the properties of extended spaces for the Einstein-spaces in General Relativity.

2 The Schwarzschild space-time and geodesic equation

The line element of standard metric of the Schwarzschild space-time in coordinate system x, θ, ϕ, t has the form

$$ds^2 = -\frac{1}{(1 - 2M/x)} dx^2 - x^2 (d\theta^2 + \sin^2 \theta d\phi^2) + (1 - 2M/x) dt^2. \quad (6)$$

The geodesic equations of this type of the metric are

$$\frac{d^2}{ds^2}x(s) - \frac{M \left(\frac{d}{ds}x(s)\right)^2}{(x-2M)x} + (-x+2M) \left(\frac{d}{ds}\theta(s)\right)^2 - \quad (7)$$

$$- (x-2M) (\sin(\theta))^2 \left(\frac{d}{ds}\phi(s)\right)^2 + \frac{(x-2M)M \left(\frac{d}{ds}t(s)\right)^2}{x^3} = 0,$$

$$\frac{d^2}{ds^2}\theta(s) + 2 \frac{\left(\frac{d}{ds}x(s)\right) \frac{d}{ds}\theta(s)}{x} - \sin(\theta) \cos(\theta) \left(\frac{d}{ds}\phi(s)\right)^2 = 0, \quad (8)$$

$$\frac{d^2}{ds^2}\phi(s) + 2 \frac{\left(\frac{d}{ds}x(s)\right) \frac{d}{ds}\phi(s)}{x} + 2 \frac{\cos(\theta) \left(\frac{d}{ds}\theta(s)\right) \frac{d}{ds}\phi(s)}{\sin(\theta)} = 0, \quad (9)$$

$$\frac{d^2}{ds^2}t(s) - 2 \frac{M \left(\frac{d}{ds}x(s)\right) \frac{d}{ds}t(s)}{x(2M-x)} = 0. \quad (10)$$

The symbols of Christoffel of the metric (6) looks as

$$\Gamma_{11}^1 = \frac{M}{x(2M-x)}, \quad \Gamma_{22}^1 = (2M-x), \quad \Gamma_{33}^1 = (2M-x) \sin^2 \theta,$$

$$\Gamma_{44}^1 = -\frac{M(2M-x)}{x^3}, \quad \Gamma_{12}^2 = \frac{1}{x}, \quad \Gamma_{33}^2 = -\sin \theta \cos \theta,$$

$$\Gamma_{23}^3 = \frac{\cos \theta}{\sin \theta}, \quad \Gamma_{14}^4 = -\frac{M}{x(2M-x)}, \quad \Gamma_{13}^3 = \frac{1}{x}.$$

The equations of geodesic (7)–(10) have the first integrals

$$\frac{d}{ds}x(s) = h \sqrt{1 - \left(1 - 2 \frac{M}{x(s)}\right) \left(h^{-2} + \frac{C^2}{(x(s))^2}\right)}, \quad (11)$$

$$\left(\frac{d}{ds}\theta(s)\right)^2 = h^2 \left(B^2 - \frac{C^2}{(\sin(\theta))^2}\right) x^{-4},$$

$$\frac{d}{ds}\phi(s) = \frac{hC}{x^2 (\sin(\theta))^2}, \quad (12)$$

$$\frac{d}{ds}t(s) = h \left(1 - 2 \frac{M}{x(s)}\right)^{-1},$$

where a dot denotes differentiation with respect to parameter s and (C, B, h) are the constants of motion.

3 The Riemann extension of the Schwarzschild metric

Now with help of the formulae (2) we construct the eight-dimensional extension of basic metric (6)

$$\begin{aligned}
ds^2 = & -\frac{2M}{x(2M-x)}Pdx^2 - \frac{2}{x}Qdx d\theta - 2(2M-x)Pd\theta^2 - \frac{2}{x}Udx d\phi + \\
& + 2\frac{M}{x(2M-x)}Vdx dt - 2\frac{\cos\theta}{\sin\theta}Ud\phi d\theta - 2((2M-x)\sin^2\theta P - \sin\theta\cos\theta Q)d\phi^2 + \\
& + 2\frac{M(2M-x)}{x^3}Pdt^2 + 2dx dP + 2d\theta dQ + 2d\phi dU + 2dt dV, \quad (13)
\end{aligned}$$

where (P, Q, U, V) are the additional coordinates of extension.

The metrics of a given type are the metrics with vanishing curvature invariants. They play an important role in general theory of Riemannian spaces. In particular the metrics for pp-waves in General Relativity belong to this class.

The eight-dimensional space in local coordinates $(x, \theta, \phi, t, P, Q, U, V)$ with this type of metric is also the Einstein space with condition on the Ricci tensor

$${}^8R_{ik} = 0.$$

The complete system of geodesic equations for the metric (??) decomposes into two groups of equations.

The first group coincides with the equations (7-10) on the coordinates (x, θ, ϕ, t) and second part forms the linear system of equations for coordinates P, Q, U, V .

They are defined as

$$\begin{aligned}
& \ddot{P} + \frac{2M}{x(x-2M)}\dot{x}\dot{P} - \frac{2}{x}\dot{\theta}\dot{Q} - \frac{2}{x}\dot{\phi}\dot{U} - \frac{2M}{x(x-2M)}t\dot{V} - \\
& - \left(\frac{2M}{x^2(x-2M)}\dot{x}^2 + \frac{(x-2M)}{x}\dot{\theta}^2 + \frac{\sin^2\theta(x-2M)}{x}\dot{\phi}^2 + \frac{2M(x-2M)}{x^4}t^2 \right) P + \\
& + \left(\frac{4}{x^2}\dot{x}\dot{\theta} - \frac{2\cos\theta}{x}\dot{\phi}^2 \right) Q + \left(\frac{4}{x^2}\dot{x}\dot{\phi} + \frac{4\cos\theta}{x\sin\theta}\dot{\theta}\dot{\phi} \right) U + \left(\frac{4M^2}{x^2(x-2M)^2}\dot{x}t \right) V = 0, \\
& \ddot{Q} + 2(x-2M)\dot{\theta}\dot{P} - \frac{2}{x}\dot{x}\dot{Q} - \frac{2\cos\theta}{\sin\theta}\dot{\phi}\dot{U} - \frac{2(x-4M)}{x}\dot{x}\dot{\theta}P + \\
& + \left(\frac{2(x-3M)}{x^2(x-2M)}\dot{x}^2 - \frac{2(x-2M)}{x}\dot{\theta}^2 - \frac{(x-4M\sin^2\theta)}{x}\dot{\phi}^2 + \frac{2M(x-2M)}{x^4}t^2 \right) Q + \\
& + \left(\frac{4\cos\theta}{x\sin\theta}\dot{x}\dot{\phi} + \frac{4\cos^2\theta}{\sin^2\theta}\dot{\theta}\dot{\phi} \right) U = 0,
\end{aligned}$$

$$\ddot{U} + 2\sin^2\theta(x-2M)\dot{\phi}\dot{P} + 2\sin\theta\cos\theta\dot{\phi}\dot{Q} - \left(\frac{2\cos\theta}{\sin\theta}\dot{\theta} + \frac{2}{x}\dot{x} \right) \dot{U} - \frac{2\sin^2\theta(x-4M)}{x}\dot{x}\dot{\phi}P -$$

$$\begin{aligned}
& - \left(\frac{4 \sin \theta \cos \theta}{x} \dot{x} \dot{\phi} + 2 \dot{\theta} \dot{\phi} \right) Q + \\
& + \left(\frac{2(x-3M)}{x^2(x-2m)} \dot{x}^2 + \frac{4 \cos \theta}{x \sin \theta} \dot{x} \dot{\theta} + \frac{2(x \cos^2 \theta + 2M \sin^2 \theta)}{x \sin^2 \theta} \dot{\theta}^2 \right) - \\
& - \left(\frac{2(x-2M \sin^2 \theta)}{x} \dot{\phi}^2 \right) U = 0, \\
& \ddot{V} - \frac{2M(x-2m)}{x^3} \dot{t} \dot{P} - \frac{2M}{x(x-2M)} \dot{x} \dot{V} + \frac{4M(x-2M)}{x^4} \dot{x} \dot{t} P + \\
& + \left(\frac{2M(2x-3M)}{x^2(x-2M)^2} \dot{x}^2 - \frac{2M}{x} \dot{\theta}^2 - \frac{2M \sin^2 \theta}{x} \dot{\phi}^2 + \frac{2M^2}{x^4} \dot{t}^2 \right) V = 0.
\end{aligned}$$

So we get the linear matrix-second order ODE for the coordinates U, V, P, Q

$$\frac{d^2 \Psi}{ds^2} + A(x, \theta, \phi, t) \frac{d\Psi}{ds} + B(x, \theta, \phi, t) \Psi = 0, \quad (14)$$

where

$$\Psi(s) = \begin{pmatrix} P(s) \\ Q(s) \\ U(s) \\ V(s) \end{pmatrix}$$

and A, B are some 4×4 matrix-functions depending from the coordinates $x^i(s)$ and their derivatives.

We shall study this system of equations at the condition $\theta = \pi/2$.

In this case we get the system for the coordinates of basic space

$$\begin{aligned}
& \frac{d^2}{ds^2} x(s) + \frac{M \left(\frac{d}{ds} x(s) \right)^2}{x(s) (2M - x(s))} + (2M - x(s)) \left(\frac{d}{ds} \phi(s) \right)^2 - \\
& - \frac{M (2M - x(s)) \left(\frac{d}{ds} t(s) \right)^2}{(x(s))^3} = 0, \\
& \frac{d^2}{ds^2} \phi(s) + 2 \frac{\left(\frac{d}{ds} x(s) \right) \frac{d}{ds} \phi(s)}{x(s)} = 0, \\
& \frac{d^2}{ds^2} t(s) - \frac{M \left(\frac{d}{ds} x(s) \right) \frac{d}{ds} t(s)}{x(s) (2M - x(s))} = 0
\end{aligned}$$

and the system of equations for the supplementary coordinates

$$\begin{aligned}
& \left(2 \frac{M \left(\frac{d}{ds} x(s) \right)^2}{(2M - x(s)) (x(s))^2} - \frac{(-2M + x(s)) \left(\frac{d}{ds} \phi(s) \right)^2}{x(s)} \right) P(s) + \\
& + \left(2 \frac{M (2M - x(s)) \left(\frac{d}{ds} t(s) \right)^2}{(x(s))^4} \right) P(s) +
\end{aligned}$$

$$\begin{aligned}
& + \frac{d^2}{ds^2} P(s) + 4 \frac{U(s) \left(\frac{d}{ds} x(s) \right) \frac{d}{ds} \phi(s)}{(x(s))^2} + 4 \frac{M^2 V(s) \left(\frac{d}{ds} x(s) \right) \frac{d}{ds} t(s)}{(x(s))^2 (2M - x(s))^2} - 2 \frac{M \left(\frac{d}{ds} x(s) \right) \frac{d}{ds} P(s)}{x(s) (2M - x(s))} \\
& \quad - 2 \frac{\left(\frac{d}{ds} \phi(s) \right) \frac{d}{ds} U(s)}{x(s)} + 2 \frac{M \left(\frac{d}{ds} t(s) \right) \frac{d}{ds} V(s)}{x(s) (2M - x(s))} = 0, \\
& \quad \left(2 \frac{(-x(s) + 3M) \left(\frac{d}{ds} x(s) \right)^2}{(2M - x(s)) (x(s))^2} - \frac{(x(s) - 4M) \left(\frac{d}{ds} \phi(s) \right)^2}{x(s)} \right) Q(s) - \\
& \quad - \left(2 \frac{M (2M - x(s)) \left(\frac{d}{ds} t(s) \right)^2}{(x(s))^4} \right) Q(s) + \frac{d^2}{ds^2} Q(s) - 2 \frac{\left(\frac{d}{ds} x(s) \right) \frac{d}{ds} Q(s)}{x(s)} = 0, \\
& \quad - 2 \frac{(x(s) - 4M) \left(\frac{d}{ds} x(s) \right) \left(\frac{d}{ds} \phi(s) \right) P(s)}{x(s)} + \\
& \quad + 2 \left(- \frac{(-2M + x(s)) \left(\frac{d}{ds} \phi(s) \right)^2}{x(s)} + \frac{(-x(s) + 3M) \left(\frac{d}{ds} x(s) \right)^2}{(2M - x(s)) (x(s))^2} \right) U(s) - \\
& \quad - \left(\frac{M (2M - x(s)) \left(\frac{d}{ds} t(s) \right)^2}{(x(s))^4} \right) U(s) + \frac{d^2}{ds^2} U(s) - 2 (2M - x(s)) \left(\frac{d}{ds} \phi(s) \right) \frac{d}{ds} P(s) - \\
& \quad - 2 \frac{\left(\frac{d}{ds} x(s) \right) \frac{d}{ds} U(s)}{x(s)} = 0, \\
& \quad \left(- 2 \frac{M (3M - 2x(s)) \left(\frac{d}{ds} x(s) \right)^2}{(x(s))^2 (2M - x(s))^2} - 2 \frac{M \left(\frac{d}{ds} \phi(s) \right)^2}{x(s)} + 2 \frac{M^2 \left(\frac{d}{ds} t(s) \right)^2}{(x(s))^4} \right) V(s) + \\
& \quad + \frac{d^2}{ds^2} V(s) - 4 \frac{MP(s) (2M - x(s)) \left(\frac{d}{ds} x(s) \right) \frac{d}{ds} t(s)}{(x(s))^4} + 2 \frac{M \left(\frac{d}{ds} x(s) \right) \frac{d}{ds} V(s)}{x(s) (2M - x(s))} + \\
& \quad + 2 \frac{M (2M - x(s)) \left(\frac{d}{ds} t(s) \right) \frac{d}{ds} P(s)}{(x(s))^3} = 0.
\end{aligned}$$

In this case the matrix A takes the form

$$A = - \begin{bmatrix} -2 \frac{\left(\frac{d}{ds} x(s) \right) M}{x(s) (-2M + x(s))} & 0 & 2 \frac{\frac{d}{ds} \phi(s)}{x(s)} & 2 \frac{M \frac{d}{ds} t(s)}{x(s) (-2M + x(s))} \\ 0 & 2 \frac{\frac{d}{ds} x(s)}{x(s)} & 0 & 0 \\ -2 (-2M + x(s)) \frac{d}{ds} \phi(s) & 0 & 2 \frac{\frac{d}{ds} x(s)}{x(s)} & 0 \\ 2 \frac{(-2M + x(s)) \left(\frac{d}{ds} t(s) \right) M}{(x(s))^3} & 0 & 0 & 2 \frac{\left(\frac{d}{ds} x(s) \right) M}{x(s) (-2M + x(s))} \end{bmatrix},$$

and matrix B has an elements

$$B_{11} =$$

$$\begin{aligned}
&= \frac{\left(\frac{d}{ds}\phi(s)\right)^2 (x(s))^4 (x(s) - 4M) + 4M^2 \left(\frac{d}{ds}\phi(s)\right)^2 (x(s))^3 + 2M \left(\frac{d}{ds}t(s)\right)^2 (x(s))^2}{(-2M + x(s))(x(s))^4} + \\
&+ \frac{2M \left(\frac{d}{ds}x(s)\right)^2 (x(s))^2 (-2M + x(s))(x(s))^4 - 8M^2 \left(\frac{d}{ds}t(s)\right)^2 x(s) + 8M^3 \left(\frac{d}{ds}t(s)\right)^2}{(-2M + x(s))(x(s))^4}, \\
B_{12} &= 0, \quad B_{13} = -4 \frac{\left(\frac{d}{ds}x(s)\right) \frac{d}{ds}\phi(s)}{(x(s))^2}, \quad B_{14} = -4 \frac{M^2 \left(\frac{d}{ds}x(s)\right) \frac{d}{ds}t(s)}{(x(s))^2 \left(4M^2 - 4x(s)M + (x(s))^2\right)}, \\
B_{21} &= 0, \quad B_{22} = \\
&\frac{-8M^3 \left(\frac{d}{ds}t(s)\right)^2 + \left(\frac{d}{ds}\phi(s)\right)^2 (x(s))^5 + 8M^2 \left(\frac{d}{ds}\phi(s)\right)^2 (x(s))^3 - 2 \left(\frac{d}{ds}x(s)\right)^2 (x(s))^3}{(-2M + x(s))(x(s))^4} \\
&- \frac{2M \left(\frac{d}{ds}t(s)\right)^2 (x(s))(x(s) + 4M) + 6M \left(\frac{d}{ds}x(s)\right)^2 (x(s))^2 - 6M \left(\frac{d}{ds}\phi(s)\right)^2 (x(s))^4}{(-2M + x(s))(x(s))^4}, \\
B_{23} &= 0, \quad B_{24} = 0, \\
B_{31} &= 2 \frac{(x(s) - 4M) \left(\frac{d}{ds}x(s)\right) \frac{d}{ds}\phi(s)}{x(s)}, \quad B_{32} = 0, \\
B_{33} &= \\
&-2 \frac{-\left(\frac{d}{ds}\phi(s)\right)^2 (x(s))^3 (x(s)^2 + 4M^2) + 4M^3 \left(\frac{d}{ds}t(s)\right)^2 - 3M \left(\frac{d}{ds}x(s)\right)^2 (x(s))^2}{(-2M + x(s))(x(s))^4} + \\
&+ \frac{\left(\frac{d}{ds}x(s)\right)^2 (x(s))^3 + 4M \left(\frac{d}{ds}\phi(s)\right)^2 (x(s))^4 - 4M^2 \left(\frac{d}{ds}t(s)\right)^2 x(s) + M \left(\frac{d}{ds}t(s)\right)^2 (x(s))^2}{(-2M + x(s))(x(s))^4}, \\
B_{34} &= 0, \quad B_{41} = -4 \frac{(-2M + x(s)) \left(\frac{d}{ds}t(s)\right) \left(\frac{d}{ds}x(s)\right) M}{(x(s))^4}, \quad B_{42} = 0, \quad B_{43} = 0, \\
B_{44} &= -2 \frac{M \left(-3M \left(\frac{d}{ds}x(s)\right)^2 (x(s))^2 - \left(\frac{d}{ds}\phi(s)\right)^2 (x(s))^5 - 4M^2 \left(\frac{d}{ds}\phi(s)\right)^2 (x(s))^3\right)}{(x(s))^4 \left(4M^2 - 4x(s)M + (x(s))^2\right)} - \\
&-2 \frac{M \left(4M \left(\frac{d}{ds}\phi(s)\right)^2 (x(s))^4 + 2 \left(\frac{d}{ds}x(s)\right)^2 (x(s))^3 + M \left(\frac{d}{ds}t(s)\right)^2 (2M - x(s))^2\right)}{(x(s))^4 \left(4M^2 - 4x(s)M + (x(s))^2\right)}.
\end{aligned}$$

Now we will integrate our system.

Remark that the equation for the coordinate $Q(s)$ is independent from others equations and can be reduced after the substitution

$$Q(s) = x(s)F(s)$$

to the equation for the function $F(s)$

$$\frac{d^2}{ds^2}F(s) + \left(\frac{M}{(x(s))^3} + 3 \frac{C^2 h^2 M}{(x(s))^5} \right) F(s) = 0. \quad (15)$$

To integrate the equations for the coordinates $P(s)$, $U(s)$, $V(s)$ we use the relation

$$\begin{aligned} \left(\frac{d}{ds}x(s) \right) P(s) + \left(\frac{d}{ds}\theta(s) \right) Q(s) + \left(\frac{d}{ds}\phi(s) \right) U(s) + \left(\frac{d}{ds}t(s) \right) V(s) - \\ -1/2 s - \mu = 0 \end{aligned} \quad (16)$$

which is consequence of the well known first integral of geodesic equations of arbitrary Riemann space

$$g_{ik} \frac{dx^i(s)}{ds} \frac{dx^k(s)}{ds} = \text{const.}$$

In our case it takes the form

$$\left(\frac{d}{ds}x(s) \right) P(s) + \left(\frac{d}{ds}\phi(s) \right) U(s) + \left(\frac{d}{ds}t(s) \right) V(s) - 1/2 s - \mu = 0.$$

Solving this equation with respect the function $V(s)$

$$V(s) = -1/2 \frac{2 \left(\frac{d}{ds}x(s) \right) P(s) + 2 \left(\frac{d}{ds}\phi(s) \right) U(s) - s - 2\mu}{\frac{d}{ds}t(s)} \quad (17)$$

and substituting this expression into the last two equations of the system we get the following two equations for coordinates $U(s)$

$$\begin{aligned} \frac{d^2}{ds^2}U(s) = -2 \frac{\left(-2Ch(x(s))^3 M + Ch(x(s))^4 \right) \frac{d}{ds}P(s)}{(x(s))^5} + \\ + 2 \sqrt{\frac{(x(s))^3 h^2 - (x(s))^3 - C^2 h^2 x(s) + 2M(x(s))^2 + 2C^2 h^2 M}{(x(s))^3 h^2}} \\ \left(h \frac{d}{ds}U(s) (x(s))^{-1} - h^2 C (x(s))^2 (4M - x(s)) P(s) (x(s))^{-5} \right) - \\ - 2 \frac{\left(3M(x(s))^2 - (x(s))^3 + (x(s))^3 h^2 - 2C^2 h^2 x(s) + 5C^2 h^2 M \right) U(s)}{(x(s))^5} \end{aligned}$$

and $P(s)$

$$\frac{d^2}{ds^2}P(s) + \frac{4Mh}{x(s)(x(s) - 2M)} \sqrt{\frac{(x(s))^3 (h^2 - 1) - C^2 h^2 x(s) + 2M(x(s))^2 + 2C^2 h^2 M}{(x(s))^3 h^2}} \times$$

$$\begin{aligned}
& \times \frac{d}{ds} P(s) = - \frac{\left(-2 (x(s))^3 Ch + 6 (x(s))^2 ChM \right) \frac{d}{ds} U(s)}{(x(s))^5 (-2M + x(s))} - \\
& - \frac{\left(2 (x(s))^3 M(1 - 2h^2) - (x(s)) h^2 C^2 (x(s) - 8M) - 6 (x(s))^2 M^2 \right) P(s)}{(x(s))^5 (-2M + x(s))} - \\
& - \frac{14 h^2 C^2 M^2 P(s)}{(x(s))^5 (-2M + x(s))} - \\
& - \sqrt{\frac{(x(s))^3 h^2 - (x(s))^3 - C^2 h^2 x(s) + 2M (x(s))^2 + 2C^2 h^2 M}{(x(s))^3 h^2}} \times \\
& \times \frac{\left(4 (x(s))^2 h^2 C - 12 x(s) h^2 CM \right) U(s)}{(x(s))^5 (-2M + x(s))} + \frac{M}{x(s) (-2M + x(s))}
\end{aligned}$$

So we have showed that every motion on orbit in usual space corresponds the motion in additional space.

Let us consider some examples.

According to ([5]) in the Schwarzschild space-time exists the cyclic orbit

$$x(s) = 6M$$

which is the solution of the geodesic equations at the condition

$$h = 1, \quad C = 3\sqrt{(2)M}.$$

In this case our system takes the form

$$\frac{d^2}{ds^2} P(s) - 1/48 \frac{\sqrt{2} \frac{d}{ds} U(s)}{M^2} - \frac{19}{864} \frac{P(s)}{M^2} - 1/24 M^{-1} = 0, \quad (18)$$

and

$$\frac{d^2}{ds^2} U(s) + 2/3 \sqrt{2} \frac{d}{ds} P(s) - \frac{1}{216} \frac{U(s)}{M^2} = 0. \quad (19)$$

The simplest solution of this system looks as

$$P(s) = -\frac{36}{19} M + A \sin \left(\frac{1}{72} \frac{\sqrt{3 + 3i\sqrt{303}s}}{M} \right) \quad (20)$$

and

$$U(s) = A \left(39 + i\sqrt{303} \right) \sqrt{2} M \cos \left(\frac{1}{72} \frac{\sqrt{3 + 3i\sqrt{303}s}}{M} \right) \frac{1}{\sqrt{3 + 3i\sqrt{303}}} \quad (21)$$

where A is arbitrary parameter.

The equation for coordinate $Q(s)$ after substitution

$$x(s) = 6M, \quad h = 1, \quad C = 3\sqrt{(2)}M$$

takes a form

$$\frac{d^2}{ds^2}Q(s) + \frac{1}{108} \frac{Q(s)}{M^2} = 0$$

and its solution is

$$Q(s) = C_1 \cos\left(1/18 \frac{\sqrt{3}s}{M}\right) + C_2 \sin\left(1/18 \frac{\sqrt{3}s}{M}\right). \quad (22)$$

At last the expression for coordinate $V(s)$ in considered case can be found from the relation (??).

It has the form

$$V(s) = -1/9 A \left(39 + i\sqrt{303}\right) \cos\left(\frac{1}{72} \frac{\sqrt{3 + 3i\sqrt{303}}s}{M}\right) \frac{1}{\sqrt{3 + 3i\sqrt{303}}} + \\ + 1/3 s + 2/3 \mu. \quad (23)$$

So the formulae (??,??,??,??) describe the relation between the properties of motion of the test particle on the orbit $x(s) = 6M$ in basic physical space with coordinates (x, θ, ϕ, t) and its map into additional space with coordinates (P, Q, U, V) .

The solution of the equation (??) relatively parameter s is

$$s = 72 M \arccos\left(\frac{U\sqrt{3 + 3i\sqrt{303}}\sqrt{2}}{2AM((39 + i\sqrt{303}))}\right) \frac{1}{\sqrt{3 + 3i\sqrt{303}}}.$$

The substitution of this value into the formulae for the coordinate $P(s)$ give us the quadric

$$95 iU^2\sqrt{303} + 2071 U^2 - 184832 A^2 M^2 + 184832 P^2 M^2 + 700416 PM^3 + 663552 M^4 = 0.$$

In the case of radial motion $C = 0$, $h = 1$ we get

$$x(s) = 1/2 2^{2/3} 3^{2/3} \sqrt[3]{M} s^{2/3}, \quad (24)$$

and

$$\frac{d}{ds}t(s) = 3 \frac{s^{2/3}}{3 s^{2/3} - 2 M^{2/3} \sqrt[3]{2} \sqrt[3]{3}}. \quad (25)$$

The system of equations for additional coordinates takes the form

$$\frac{d^2}{ds^2}P(s) + 4 M^2 \sqrt{2} \frac{d}{ds}P(s) (x(s))^{-2} \frac{1}{\sqrt{\frac{M}{x(s)}}} (-2 M + x(s))^{-1} -$$

$$\begin{aligned}
& - \left(2x(s)\sqrt{\frac{M}{x(s)}} + 6M\sqrt{\frac{M}{x(s)}} \right) MP(s) \frac{1}{\sqrt{\frac{M}{x(s)}}} (-2M + x(s))^{-1} (x(s))^{-3} - \\
& \quad - \frac{M}{x(s)(-2M + x(s))} = 0, \\
& \quad \frac{d^2}{ds^2}Q(s) - 2 \frac{\left(\left(\frac{d}{ds}x(s) \right) (x(s))^4 - 2M \left(\frac{d}{ds}x(s) \right) (x(s))^3 \right) \frac{d}{ds}Q(s)}{(-2M + x(s))(x(s))^4} - \\
& -2 \frac{\left(-M \left(\frac{d}{ds}t(s) \right)^2 (x(s) - 2M)^2 - \left(\frac{d}{ds}x(s) \right)^2 (x(s))^2 (x(s) - 3M) \right) Q(s)}{(-2M + x(s))(x(s))^4} = 0, \\
& \quad \frac{d^2}{ds^2}U(s) + 4/3 \frac{U(s)}{s^2} - 4/3 \frac{\frac{d}{ds}U(s)}{s} = 0,
\end{aligned}$$

and

$$V(s) + 1/2 \frac{2 \left(\frac{d}{ds}x(s) \right) P(s) - s - 2\mu}{\frac{d}{ds}t(s)} = 0.$$

After substitution here the relations (??,??) we find the solutions

$$\begin{aligned}
P(s) &= 1/2 \frac{\sqrt[3]{s}\sqrt[3]{2}\sqrt[3]{3} (3C_5 + s\sqrt[3]{2}\sqrt[3]{3}M^{2/3})}{-3\sqrt[3]{M}s^{2/3} + 2M\sqrt[3]{2}\sqrt[3]{3}}, \\
Q(s) &= C_3 s + C_4 s^{4/3}, \\
U(s) &= C_1 s + C_2 s^{4/3},
\end{aligned}$$

and

$$V(s) = s/2 + \frac{C_5}{s^{2/3}}.$$

The linear system of geodesics for additional coordinates (??) may be used for the studying of the properties of a basic space. In particular the sequence of the matrixes

$$E(s), \quad E_{;s}, \quad E_{;ss}, \dots,$$

where

$$E_s = \frac{dE(s)}{ds} + \frac{1}{2} [A(s), E(s)]$$

and their invariants are important characteristic of a basic space.

Remark that for a given example the matrix $E(s)$ has a property

$$\text{Det}(E(s)) = 0, \quad \text{Trace}(E(s)) = 0.$$

More detail consideration leads to conclusion that in general case for the matrix $E(s)$ the condition

$$\text{Trace}(E(s)) = R_{ij}\dot{x}^i\dot{x}^j$$

is obeyed, where R_{ij} is the Ricci tensor of the basic space.

The generalization and the interpretation of these results will be done later.

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